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PUERTO RICO NUCLEAR CENTER

RADIATION DAMAGE IN ORGANIC CRYSTALS

Progress Summary Report No. 3

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STUDY OF RADIATION DAMAGE IN ORGANIC CRYSTALS USING ELECTRICAL

conpuctrvity

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Progress Report #3

Work performed at Puerto Rico Nuclear Center

Rio Piedras, P.7., under U.S. Atomic Fnergy

Comission Contract AT(40~i)-i833 (Project 14)

Senvary 1965

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This project is concerned with the effects of radiation on organic crystals. It is felt that such studies on well defined crystal-

Line

structures can provide a firm foundation for a later study of more complex materials including those of direct biological interest. We have

chosen anthracene as the initial material for study because this subst

has been studied more than any other organic material,

The effect of neutron irradiation on anthracene has been studied

Previously by Kamandeur (!+?), but to the best of our knowledge, no other

work on this subject has appeared since then. Since Komandeur's work was

done very early in the history of organic conductivity, we felt that it

would be valuable to reopen and expand this work to include more recent

developments such as the introduction of charge-injecting electrodes (),

and the application of space-charge-limited current theory to molecular

chemistry (1,2,3,4,5,6,7,8,9,10,11,12)

Below 19 « summary of the results obtained during the period

January 1964 to December 1964,

Our preliminary results of neutron radiation damage presented

4m our Progress Reports #1 and 42 indicated that in electrical conductivity

measurements a change in the

duration current occurs after a very low

dose of radiation. These measurements were made using « Kalinann-Popa(!?)

cell with Na₂SO₄ solution as one of the electrodes and Na₂S₂O₃ solution

% 2

the hole injecting electrode. Subsequent measurements with a more sensitive

electrode system recently developed at our laboratory indicate that the

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changes detected are most likely an electrode effect rather than &

change in the crystal itself. However, our most recent work shows that after irradiating with gamma rays or x-rays and making the

with the improved electrode system a definite change to the electrical conductivity of the crystals was detected while no change was observed after neutron irradiation of doses comparable to those reported previously, These results comprise the main part

of this report.

SECTION 2. Crystals

As discussed in Summary Report #2 we have used an improved Kattmann-Pope() technique to grow large anthracene crystals from

solution. The main features of this technique is to use a solvent comprised of equal parts by volume of dichloroethane, cis-dichloroethylene and trans-dichloroethylene and to carefully control the

rate of cooling of the solution. A complete description of this

technique will be submitted for publication in the near future.

Using this technique we have grown all the crystals needed

for our experiment. These crystals,

(approximately 2 cm²), and appropriate thickness which ranges from

20 to 100 microns, were especially suited for our measurements.

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Crystals were irradiated using the following radiation sources

(a) Pu-Be neutron source

(b) the thermal column at the PRNC research reactor

(c) 4 Co-60 gamma ce

(3) © 120 kV x-ray radiotherapy unit

Our Pu-Be neutron source in 10 curie source imbedded in paraffin.

The crystals were placed right next to the source where the flux is approximately 2×10^5 n/cm² sec and the average neutron energy is approximately 2 Mev. Assuming the neutron absorption coefficient of anthracene to be approximately equal to that of other organic substances the neutron radiation absorbed by anthracene is approximately 3×10^9 rad/n/cm² (!5) where 1 rad equals 100 ergs/gm. Due to the fact that our Pu-Be neutron

source has such a small flux the doses given in reasonably short times

(3 days to 1 week) were of the order of a few ra

The crystals irradiated with the reactor were irradiated with thermal neutrons (0.025 eV). The flux used was approximately 10^9 n/cm²/sec and the doses given ranged from 1000 to 12,000 rads approximately.

The Co-60 gamma cell used Co-60 irradiation

of our crystals that

output

of 780 r/min. and the energies of the Co-60 photons are 1.33 and 1.17 MeV.

By measuring the absorption coefficient of anthracene for the Co-60 photons,

which turned out to be 0.118 cm⁻¹, we calculated that the energy absorbed by

anthracene is 293 erg/gal. The doses given ranged from 8×10^2 to 10^4 r,

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The x.

y unit used to irradiate our crystals w

2120 kwp

radlotherapy unit with an output of 325 r/min, By measuring the absorption

coefficient of anthracene for this radiation, which turned out to be

1

0.671 en", we calculated that the energy absorbed by anthracene is

rl. the doses given ranged from 8×10^2 x to 10^8 x,

296 erg/gm

SECTION ITE. EXPERIMENTAL ser UP

?The experinent

set-up {6 shown tn Fig

?The electro

configuration system chat ve have developed and used ia as follows,

?The non

injecting electrode consiate of a plece of transparent

conducting glass (evaporated tin oxide on one of the surfaces) on

which 4 small drop of 1M Na₂S₂O₈ solution is placed. The crystal is then layered on top of the Na₂S₂O₈ solution drop which then spreads. Producing a very good electrical contact between the crystal and the conducting gaskets. The injecting electrode on the opposite side of the

crystal consists of a drop of Net-1 solution

on which a platinum

wire is attached

is held in place by 9 simple splayed supports. In

order that the reproducibility of the saturation value of the current be within a factor of two, the concentration of the T_p in the Na₂S₂O₈-T_p solution has to be within well defined limits. This concentration is

achieved by diluting «

saturated solution of Lugine in IM NaI at room

temperature to 25-50% dilution, with a saturated solution the

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-s-

Reproducibility of the saturation current values is very poor.

Voltages between 2 and 500 volts were applied to the crystals

from a Keithley Model 240 regulated power supply. The current

measured using a Keithley Model 600A electrometer.

The illumination of the injecting electrode was achieved

using the light from mercury lamp, (Osram #80 100 W/2) Tais Light

was passed through either a Corning Color filter #S 5-74, (λ_{max} = 6360Å)

or through a Corning Cut-off filter #CS 3-72 which cuts-off at 4300Å* or

through a Corning 3-70 which cuts-off at 490Å". As is shown in Fig. 1 the

path of the light, after leaving the filter, through the conducting

and crystal and then into the NaI solution.

SECTION IV. EXPERIMENTAL RESULTS

Typical current-voltage characteristic curves on a log $\{v_e$

og V plot are shown in Fig.

2. The changing parameter in these curves

is the concentration of the iodine in the NaI solution. The values

of each curve are the average values gotten from many successive measure-

ments on a crystal. Measurements taken on large number of similar

crystals give similar results. In all the measurements represented in

Fig. 2 the electrode

was not illuminated. From measurements of this

type it was found that the iodine concentration most convenient for our

measurements is a 25% dilution.

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No detectable changes were obtained after irradiating the

crystals with the sources 12? Leatee 4m fa) and () of Section 1f

ing the doses indicated in this section.

Fig. 3 shows the space charge field current vs log

va Log V plot with one window 42um. The Imjccins electrode. The

Aodine concentration of ch: ?nfactine aleervude for al? our measurements

was 25% of the saturation level, Jo uhan a9 Pp. 2 the rapraduethley

te good using this concentration. ?urve i af Fig. 3 show the results

beained when the crystal was not (Jluninatec. Curve #1 shovs the

Fesulte obtained when the injecting electrode was i:lumtnated through

the Corning Color ftlter #3 5-74, which his « maximum exenamteaion at

D. = 4360". curves #1tr and SIV show the results obeatned when che

sngecting electrode was {Hluntnaced through che Corning Cut-off f1tere

fs 9-72, (out off ax = atoan'), and 4s 3-70 (cut off at A a4900%*y,

respectively. The relative dlspacenents of these four curves ace

explained by the fact ehae the Light Lo parctally absorbed by the crystal

and part by the injecting odine electred

Av previously stated che injecting electrode was {1luminated

with light from a high peeseure lig source,* and

shown in Fig. 1 pa

frat through the color or cut-off filter then through the crystal and then is absorbed by the electrode. The displacement of curve IT of Fig. 3,

4m which the color filter was used, from curve I, taken in the dark, in the

* tmcenaiey with filter fe9 5-74 170pure/om?

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a

region below the saturation current, is explained by the fact that part

of the light transmitted by this filter is absorbed by the anthracene

crystal. This results in a redistribution between the free and trapped

carriers producing higher current values for the same voltages. The

higher saturation values in the {itwminated cenes,(ts due co the

orption of the li

in the iodine solution near the crystal surface

Anereasing the concentration of the disassociated todine.

In curve TIT, taking account of the absorption coefficient of the anthracene and the transnission of the cut-off filter used, lets Light te absorbed tn the crystal than in curve IZ resulting in a analler dtsplecoment compared to the dark current curve,

However, the

uration value of the current of curve III is greater than that of curve IT because more Light is absorbed in the todine

Since the cut-off filter, CS 3-72, transmits 411 the long wavelengths of the source which are absorbed in the iodine, and also there is less

Light absorbed in the ant!

ene. The cut-off filter used for curve IV

cuts off at a longer wavelength than that used for curve III, and taking

into account the spectral distribution of the source explains the

relative position of this curve

Fig. 4 shows the steady state space charged Limited dark

current-voltage curves for a crystal before and after gamma irradiation.

The parameter for these curves

is the time of exposure to the gamma ray

source. It is seen that, after irradiation, the slope of the curve at low

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low current values is smaller than at higher current values. For high
Arrattation Jones this slope at low current values is two. This square

law extends to larger

ages the longer the time the crystal is

irradiated

is seen, it is also seen that the saturation current, which is

Almost the same in all the curves, occurs at higher voltages when
the irradiation time is longer. The curves in Fig. 5 were taken on
a crystal which was irradiated with gamma rays for 3 hours. Curve 1
is the dark current and curves IT, ITI and IV were taken under 100%
illumination of intensity of 25, 50 and 100% respectively.

Fig. 6 shows the x

results of measurements taken on a crystal

after irradiating with x-rays. As seen the curves of Fig. 6 are

similar to those of Fig. 4.

SECTION V. THEORY OF SPACE CHARGE LIMITED CURRENTS

It is well known that defects of solids produce changes in

the optical and electrical properties of solids. Thus the measurement

of these properties yields information about these defect states. To

order to make practical electrical measurements in insulators one must

enhance the free carrier density. One way of doing this is by injecting

free carriers into the insulator. The current will then be limited by

the space charge injected into the crystal. The steady state current

for ideal crystals discussed by Mott and Gurney, (16) depends on

the square of the applied voltage and is inversely proportional to the

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cube of the crystal thickness. The theory of the steady state space

charge limited current in real crystals was given by Rose(*) and Lamperc?),

From the literature: of the SCL in crystals with traps and dislocations

energy level, the trap density and depth can be computed. The theory

the ©

space charge Limited current for the case of an ideal

crystal was first given by Many, Stehony, et al and Levinson? and

independently by Helfrich and Markl?) the theory of the craneiont

United current for x

1 exystale vas given dy Many and

From transient measurements of the SCLC it is possible to compute the mobility and the trapping time of the carriers. These results combined with those of the steady state enables one to compute

the capture cross

section of the traps. Hence, the measurement of the

space charge limited current in insulators is a powerful tool for determining the characteristics of defects. Since in this report the

Results of steady-state space charge Limit

Given, a brief review of the theory for the steady-state function presented

below,

Following the derivation, the equations governing ϕ :

steady state current flow in a one dimensional plane geometry crystal,

are the following

1. $\nabla \cdot \mathbf{j} = 0$

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Bq. 2

were 70) and U_{ae} in quasi-steady state equilibrium.

The J is the current density; E is the electric field intensity; ee is the magnitude of the electric charge; μ is the electronic mobility; ϵ_0 is the static dielectric constant of the insulator; D is the diffusion constant for electrons; n_0 and $n_0(t)$ are the densities of the free and trapped electrons respectively; and W and τ are the values of W and τ respectively in the bulk neutral crystal in thermal and electrical equilibrium, (at 48, 20 applied voltage).

In solving these equations the diffusion term in eq. 1 is neglected, which is reasonable for voltages greater than $AZ \approx 0.025$ volts. The solution to these equations for the ideal crystal, (no traps), with the boundary condition $E=0$ at

was given by Mott and Gurney by substituting the expression for n in Bq. 2 into Bq. 1 and integrating with the above boundary condition. The result for J is

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Tampert gave the exact solution for 4 erysta] containing traps

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the trap level energy as measured from the botton of the conduction band

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and OF te the recat:

vel of the crystal at the anode. In this case the

current is diminished by a factor @ compared to the ideal case, (Bq. 3),

cece AE ta the magne ofthe crap depth, andy andy are te
duaeenn of mflable sates tn the coast on aad and rep regetéon

im the current occurs and the value given by 8.3 for the ideal case ts

reached.

Lampert shoved that

current-voltage characteristics curves

ween plotted on a top T versus tog V plot wilt be confned to lte in «

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erlangle as shown in Fig. 7. Por voltages lesa chan Vi the current

Will follow Ohm's Lav bectuse the cumber of injected carriers $n(x)$ is 1

chan FI , che mumber of Gerrans t=

absence of injection. For

a-0

voltages greater than V'' in a trap free crystal, the current will obey Child's Law as given by Bq. 3. In the case of # real crystal the departure from Oha's Law will occur at # higher voltage due to the fact

that only & fraction of the total inject

carriers ts free. I? all the

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+12.

traps are at a single energy level the current vill again follow

av? ay but will be below the child's Law by a factor O

TEL

by Bq. 5. At the voltage, WV, when the number of injected carriers

given

4s approximately equal to the number of trap

in the crystal, a large

rise in the current occurs for small increases in voltage until the
child's Law value is reached. For higher voltages, since all the

traps are filled, the injected carriers

in an ideal crystal and

child's Law is followed.

For crystals which contain trap levels that are distributed

in energy, ϕ will no longer be a constant but rather a function of

the voltage. In these cases, when the current departs from Ohm's Law

it will follow a voltage dependence $J \sim (U - U_0)^2$. One can see that

from a simple steady state current-voltage measurement made on an

insulator with an injecting electrode a large amount of information

can be ascertained. By taking the ratio of the experimental value

to the ideal Child's Law value for the current J is evaluated.

From the voltage dependence of J the energy distribution of the

carriers

trap levels can be determined. From the J vs U curve the trap density,

N_t can be evaluated, in the case where J is proportional to $(U - U_0)^2$ of the

voltage the energy depth of the traps can be found by using Eq. 5

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SECTION VE. concuusroNs

As {s seen im Fig. & the current-voltage curve is effected by the gamma radiation in such a vay that for « given voltage the current value is lovered. The larger the radiation dose, that te, going fram curves TT to V, che lover 1s the current, There 1s also a change in the slope. Beginning with curve IT the current at lov voltages departe from av^n dependence, where $n > 2$, to v^2 dependence. In curve Z + V eependence holds t111 the breakdom voltage. For Bigher radiation doses the current are lover chan in curve F but the v^2 dependence still holds.

These curves are interpreted in terms of che space charge Lintted

theory. Results simtlar to curve T, the measurements made on the thin

anthracene crystal before irradiation, have been reported by Mark and

nergeien(®). rhe

curves can be explained in terms of an exponential trap distribution. The lowering of the current value for a given voltage going from curves IT to Y, are interpreted in terms of traps which a

created by the radiation. As N_t , the density of traps increases

a, the

number of free carriers at a given voltage is decreased. The gradual

change in the slope to a value of 2 for higher radiation doses strongly indicates that the traps introduced by the radiation are all at the same discrete energy level. Using the equations of the space charge limited

current theory outlined in Section V, the parameters of these traps can

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- ue

ϕ calculated. From curve IV of Fig. 4, the value of the voltage where the curve departs from a square law dependence is seen to be about 200 volts.

Taking this to be the value of the trapped filled limit voltage for the

aLh

TFL.

Introducing aLh and using this value in the equation, $V_e = LEA$

where L is the crystal thickness, J the elementary charge, and ϵ the

dielectric constant, N_t , the trap density, is calculated to be $2 \times 10^{13}/\epsilon$.

From the ratio of the current of curve IV and the ideal V_e curve, the ratio,

8.18×10^{-4}

of states in the conduction band can be $10^{14}/\epsilon$, the value of AL , the

Using Eq. 5 and, $N_t = 10^{13}/\epsilon$ and taking M_e , the

umber

trap depth, 16 found to be 0.92 eV.

From curve 1 of Fig. 4 it is seen that a saturation value of the current is reached. As is seen from curves IT and TET this saturation

value is not affected by the radiation. However, the voltage at which

saturation occurs depends on the radiation dose, that is on the number of

traps introduced. The reason for this is that as the number of traps is

increased the number of free carriers is decreased at any given voltage.

The electrode injection, which at the steady state is dependent on the

number of free carriers

WINL in the case of greater trap density, hold

for higher voltages, For the higher radiation doses, curves TV and V,
4 breakdown occurs before the current saturation is achieved. In Fig. 5
the Light intensity dependence of the current is shown for a crystal
where « large number of traps has been introduced. curve T of Fig. 5
was taken with a crystal in the dark, and curve TT was taken when the

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2s

exposed was illuminated by light which passed through filter C5 3-70,
(cut-off at 4500Å). Taking this HeNe intensity as 100%, curves ITE
and 1V were measured at intensities of 50% and 25% respectively.

The results taken on a crystal exposed to 120 kVp x-ray
as shown in Fig. 6, The time of irradiation differs for each curve
and was adjusted so that the total amount of energy absorbed in the
crystal was the same as the equivalent curves for the gamma radiation
shown in Fig. 4. It is seen that qualitatively and quantitatively the
curves are similar.

our work so far shows that anthracene crystals are damaged

by radiation from gamma and x-rays and that this damage can be detected:

by electrical conductivity measurements, This type of measurement can be used to measure the magnitude of the damage but does not provide information about the nature of the damage. we hope that the experiments we have planned for the future will enable us to obtain information

concerning the nature of the damage

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?Typical park current-Voltage characterstatie

Curves for Different Tordine Concentrations

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T- ELECTRODE USED 13 4 SATURATED

[> SOLUTION OF IODINE IN H₂O

I- ELECTRODE 15 4 SATURATED

SOLUTION DILUTED TO 25%.

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CURRENTS CAMPS)

Fig. 3, current-voltage curve for @

and under illumination with

crystal in the Dark

Various Filters

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CORNING FILTER cog-7H

Pomel 3? 83-72

coon » C3370

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CURRENT CAMPS)

Fig. 4.

Derk Current-Voltages Curves of a Crystal
before and after Exposure to Gaama rays
for Different tines

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ENERGY ABSORBED /W THE CRYSTAL QUE
gL TO GAMMA /RRADATION

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CURRENT (AMPS)

Fig. 5. cures

c-Voltage Curves

8 Function of Light

Intensity for a Crystal which was Exposed to

High Gamma ray Radiation Dose.

6

Intensity

ENERGY ABSORBED IN CRYSTAL DUE

TO 70 KEV GAMMA RADIATION IS 2.7410⁻¹⁸ ERG

PER DARK

AT 25% LEUCRYSTAL INTENSITY

AT 50% LEUCRYSTAL INTENSITY

AT 100% LEUCRYSTAL INTENSITY

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Fig. 6. dark Current-Voltage Curves of « Crys!
and after exposure to woraye fer DLE

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ENERGY ABSORBED INV CRYSTAL DUE

TO A-RAY IRRADIATION

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