

# PRNC164

PANC-164

PUERTO RICO NUCLEAR CENTER

GAS STOPPING POWER MEASUREMENTS FOR ALPHA PARTICLES

Eddie Ortiz and Gilberto M. Arenas Rositlo

May 1973

(OPERATED BY UNIVERSITY OF PUERTO RICO UNDER CONTRACT  
NO. AT (40-1)-1833 FOR US ATOMIC ENERGY COMMISSION

---Page Break---

University of Puerto Rico

Mayaguez Camps

GAS STOPPING POWER MEASUREMENTS FOR ALPHA PARTICLES

y

Gilberto M. Arenas Rositlo

A thesis submitted in partial fulfillment  
of the requirements for the degree of

Master of Science

(Nuclear Engineering)

May 1973

Measurement of Dopamine ~ 7

Director, Graduate Studies Date

---Page Break---

A method has been developed to measure the stopping powers of  
air for alpha particles using a  $^{28}\text{Si}$  semiconductor detector mounted  
opposite a natural  $\text{Am}^{241}$  alpha source in a gas chamber and a 102-channel  
data acquisition system.

By varying the gas pressure and taking an energy loss measurement  
at each pressure, the stopping powers and molecular stopping cross  
sections were calculated; also, range-energy relationships were measured  
simultaneously.

Terrestrial data are given for air, plus O<sub>2</sub>, CO<sub>2</sub>, N<sub>2</sub>, O<sub>3</sub>, Ar, Kr

From the Gil, Onli, » Col, Ce and OH, for alpha particle in the

range 0.63-5.0

efficiencies range from approximately 4-5 per cent at the highest.

energies to approximately 9-10 per cent at the lowest energies:

of alpha energy. The

introduced probable error in the

---Page Break---

ACRIFORMLEICIDENT

It is the author's wish to express his sincere gratitude to

those who helped him to complete this project

To IR, BODIE ORTIZ, Senior Scientist, Nuclear Engineering

Division, Puerto Rico Nuclear Center, for suggesting the problem,

for his dedicated guidance and competent supervision of the work.

To MR. AVIVA B. CATALAN, Professor, Nuclear Engineering  
Department, UPDAR, for her constructive criticism and for her  
valuable advice.

To MR. DONALD S. SASSCER, Head of the Nuclear Engineering  
Department, UPDAR, for his sincere interest in the author's work.

To the Colostan Institute of Mines and Petroleum-ICSTEE  
Foundation for granting the financial assistance needed for comple-  
tion of this project

to the personnel of the Library and Reproduction office of  
the PRIC for their exemplary cooperation,

---Page Break---

penitron

?To my dear Esther

and Alejandro

WV

---Page Break---

TASLE OF coursrrs

List oF maiz

List oF rows

CRAPTER I = meTRoDUCTTCY

CHAPTER IT = REVIEY oF 11

CHAPTER IIT ~ THE ITEUCTION OF ALPHA RADIATION WITH MATTER

RATE

3-1 GoLLision toes and Stopping Pover

3.2 Stopping Cross Section

3.3 Bethe!

Gurr Ty - ppEmen,

Theory

SEIT AID PROCEDURE

41 Detection Syetea Daseription

402 Brergy Lona Nessurszent

4-3 Cross Section Calculations

GuPIR V = RESULTS WD covetastoxs

BIBLIOGRI

aro

---Page Break---

Table No.

bt

hd

43

bt

at

a2

43

ws

45

46

rr

ry

rt)

410

Some chemical and physical properties of gases used

in this work

Range-energy relationship of alpha and hydrocarbon

Gases for alpha particles

Molecular stopping cross sections of air and hydro

carbon gases for alpha particles

Comparison of molecular cross section results with

data found in Literature

Instrumentation Equipment settings for experiments

Range and Stopping Power Experimental Data for 5.477

alpha particles in air

Range and Stopping Power Experimental Data for 5.477-eV

alpha particles in C:



Range and Energy Loss

alpha particles in 02%

Printed Data for 5.477 MeV

Range and Energy Loss Experimental Data for 5.477-127

alpha particles in G33

Range and Energy Loss Experimental Data for 5.477-MeV

alpha particles in Cai,

Range and Energy Loss Experimental Data for 5-47

alpha particles in 2335

Range and Energy Loss Export

alpha particles in 0

wental Data for 5.477%-NeV

Range and Sherry Loss Sperinental Date for 5.477=loV

?alpha particles in Op

ge and herzy Loss Exporisontal Date for 5-47-tfe7

adpha particles in ig

6%

66,

67

---Page Break---

LIST OF TABLES coins

Ant

bi?

13

para

Range and Energy Loss Experimental Data for 5.47 MeV

?alpha particles in A

Range and Energy Loss Experimental Data for 5.47 MeV

alpha particles in C2,

Range and Energy Loss Experimental Data for 5.47 MeV

?alpha particles in ice

Range and Energy Loss Experimental Data for S-L71-UeY

alpha particles in Freon-14

Page

---Page Break---

Last of FICuRES

Figure to.

401? Schematic diagram of the experimental arrangements

42 Beportzontal set up for stopping power and range

4.3 Gloss up picture of the experimental set up.

Photo of chamber Alpha instrumentation.

45 Calibration curve for the stopping power measuring system.

46 Several energy spectra of alpha particles after they

eroizy 5. e59 of air at various goo proscires and

at 22,

47 Resainins cnorsy of 5.t7atior i247 atpha particles

after traversing a cartain poth Xytpy corresponding

to tio chasber gas procure for aif" ang hydrocarboa,

seo,

48 Ranaining enorey of 5.L7-Hev in" atphe particles

after travaroiag a cortain path Xptny corresponding

?to tho chasber gas pressure for Ay lity 02 and Coe

469 Resaining onensy of SI7-tier 2:4! atona particles after

teavercins a certain path Xgen, cormasyoncing to the

ehasber gaa proacure for Kx)? ffo snd Freon te

4-10 Brorey lose of alpha particies tn atzy Oslig end Ogee

411 Rrarcy eas of alpha particles tn Gig, CR and C6

412 Berry Loss of alpha partictes tn A, Op and OOze

499 Bharzy 2000 of atzhn particles An tzy Kezy 10 and

08 lhe

sth Ranzo curves for tho tn" athe partieloa (5447 tov

anit

?weep:

otis cicrcy) and for alcla particles after

ng very thin Gx and TA foila tn air.

vat

6

"

8

8

mn

32

35

---Page Break---

417

ete

419

An

beat

hae

423

be,

3 ator

2 in latiane (GI)

Range curves for tho 1x1 esis purtdeles

Anitiel ?otic encray) end for elvis pert

?woaparsing vory thin Ou and Mi foils in Zt

ticles (5.47 Mev

© pareicica after

respecting very thin Gu and M foils in Frozane (Osi)

na petioles (5.47 Nev

Range curves for the 1% tdeles (5047 Mov

Aitéel ?netic encr7) and for alte particles after

frepessing very tin Ouand Ti foils in Propylene

(3).

Rengoweucy relatica of elow elpha particles in aie

and hydrocarbon gasca at standard conditions.



Range curves for As<sup>238</sup> alpha particles (5447 Yev) in  
various materials (graphite, air, etc.)

Range curves for Am<sup>241</sup> alpha particles (5447 Nev) in  
CO<sub>2</sub> gas and in air

sections of air, Oni and

Molecular stopping cross sections of N<sub>2</sub>, O<sub>2</sub>, and  
O<sub>3</sub> for alpha particles. \*

Range-energy relation of low alpha particles in air  
at standard conditions

Page

a

ar

33

---Page Break---

Guprr 1

perromvorro

In recent yours there hes

8 considereble interest tn date

Telating to stopping pover of chargod partistes dn various saterials,

The aigniticans of this paraneter aay bo rartioulary appreciated by  
dleaentary-particls piystctote and nuclear plystetate in view of the

fact that the prectaioa vith witch nuclear reection cross sections can

de measured often depends on the

section of the target =

accuracy with which the stopping cross

section is known. Health physicists need to stop»

Physical power considerations for radiation protection purposes because bio-

logical, chemical and physical effects produced by the charged particle

are deposited in a cell like human body tissue (of composition

10613, 0.124%, 1.4% and 0.73.

by weight) depend among other

factors on the absorbed dose and on the Linear Energy Transfer of the

charged particle involved,

Material composition in the material through which the radiation

4a passing (stopping material), 19 closely related to the energy loss  
by the penetrating radiation, The o

of the penetrating

charged particle per unit path length in the stopping material is  
called the stopping power of the material. Stopping power depends  
upon the various mechanisms in which radiation interacts with the  
individual atoms and molecules and is an expression of the average energy  
loss of a large number of individual interactions

---Page Break---

Stopping power measurements have been carried out in various  
solids, liquid and gases as stopping materials. In this work, a  
method is developed for measuring the stopping power of several or-

air ionization detector. By varying the gas pressure and taking  
energy loss measurements at each pressure, stopping power curves  
and molecular stopping cross sections as a function of alpha energy  
have been calculated. Also, range curves and range-energy relations  
are given for 0.3 ~ 5.4 MeV alpha particles in air and hydrocarbon gases

were developed simultaneously. Results agreed within a 5 = 10% of accuracy with respect to literature related data

---Page Break---

As indicated by recent professional literature, there is a renewed interest in the variation of stopping power of cases for heavy charged particles (with  $\ll$  mass very much greater than the mass of the electron).

A large amount of experimental work has been done and theoretical expressions for stopping power have been developed showing various degrees of accuracy.

R.B.J. PAIR (11) obtained Linear Energy Transfer (LET) curves for 1 to 8 MeV alpha particles in hydrocarbon gases and hydrogens

E. ROTONDT (13) measured stopping powers for 0.1 to 5.3 alpha particles in N<sub>2</sub>, O<sub>2</sub>, CO<sub>2</sub>, and CO, by differentiating the range-energy curves of those gases obtained by means of a variable-pressure gas cell in which alpha particles lose part of their energy and were detected by a semiconductor detector. These measurements have quoted accuracies of 8 per cent between 0.1 and 0.5 MeV, 5 per cent at 1 MeV and 3 per cent at higher energies,

PDs BOURLAND, WoK. CHU and D. POWERS (4) measured stopping

ero

etions for elphe particles in a differentially pushed gasncall

system froa 300 Key to 2

7 £m Hoy Higy Og Mlgy N0, CD, Czy Cty

Calley Callgs Callés O38 and (Gig).

GoD. KERR, Lat, RATER, Me UY

ROD and A,

1. WALTER (8) roported

experinatel data for air, tig, a, Kr, OW, for alphs particles

---Page Break---

4m the energy ranges X00 Kev ~ 5 lv. Molecular cross sections of the

above cases were calculated within an accuracy of approximately 4 per cent of probable error at the highest energies to approximately 12 per cent at the lowest energies.

PeJe WALIH (15) measured molecular stopping cross

sections for 0.3 to 5 MeV alpha particles within a 10 per

cent of accuracy

at the lowest energies, using a variable-pressure gas cell and a constant separation distance between the alpha source and detector.

WP, JESS and J. SADAUSSIS (7) determined range-energy curves

in the region 0.5 to 5 MeV for alpha particles with a collimating absorber

of known thickness and an ionization chamber. The ionization change from a

single alpha particle is collected in the ion chamber and amplified

by means of the vibrating-reed electrometers

/A, BETHE (3) reported range-energy relations for 0 to 6 Mev  
alphas particles and developed a theoretical treatment on stopping  
power for heavy charged particles ( See Section 3.3 )+

---Page Break---

curr m1

THE INTERACTION

OF ALPHA RADIATION WITH MATTER

ALL radiation measurements depend on the interaction of the  
radiation with matter, The nature of these interactions therefore,  
forms the basis of a discussion of the characteristics thereof, and  
an outline of the principal interactions of charged particles like  
alphas is developed in this section.

The primary mechanism of energy loss in penetrating charged  
particles is due to Coulomb interactions between atomic electrons and  
electrons with the charged particles; to a lesser extent elastic and  
inelastic collisions with electrons and nuclei are also of significance.  
If the energy transferred to an electron is only enough to



raise it to a higher level in the atom, the process:

4s called excited~

tion; if the electron 4s given enough energy to separate it completely from the atom, the process is called ionization. the two

process

are closely associated and together they constitute "energy

loss by collision". At energies several times the rest energy of the moving charged particle excitation and ionization account for the major part of the energy loss in all materials. Some of the electrons ejected in ionization process

have enough energies to produce further

ionization themselves; such electrons are called secondary electrons

---Page Break---

If other phenomena occur when alpha particles traverse

matter. Interactions with the Coulomb fields of atoms, and particularly of atomic nuclei, result in a change in the direction of motion of the heavy particle (the "heavy particle" shall henceforth refer to a particle of a mass greater than that of the electron). Inelastic collisions with electrons are by far the most important processes by

Which @ penetrating !

any charged particle loses its kinetic energy

when the velocity of the particle,  $v$ , is much greater than the velocity

of the electron (in hydrogen  $v_0 = 2.183 \times 10^8$  cm/sec, while for

$4 \times 10^9$  cm/sec for  $^{238}\text{U}$  alpha,  $v = 1.623 \times 10^{10}$  cm/sec). If kinetic energy is

conserved, the process is called elastic scattering; such scattering

is of minor importance for heavy particles but of great importance for

electrons

Scattering through a large angle entails a large acceleration

of the charged particle. This in turn may result in the emission of a quantum of electromagnetic radiation, known as bremsstrahlung, which becomes important when electrons are involved.

31. Larmor Loss and Radiating Power. According to the classical point of view, a moving charged particle loses energy to an electron

by imparting to it an impulse proportional to the strength of the Coulomb force and to the time during which this force acts. Since

the momentum acquired by the electron is proportional to the time

during which the interaction takes place, it is inversely propor-

portional to the velocity  $v$ , of the moving particle. The energy acquired

---Page Break---

by the electron, and hence the energy lost by the particle, must

therefore, be proportional to  $1/v$ . This on the classical picture  
of energy loss shows  
that it is proportional to the electron density in the medium and inversely  
proportional to the energy loss per unit path length (sp

proportional to the square of the velocity of the particle, hence  
the penetrating particle is moving slowly ( $v \ll v_g$ ) then the average  
energy loss per unit length of path to electrons decreases to zero, proportional to the  
velocity of the penetrating particle if the average energy loss  
due to elastic nuclear collisions is increasing proportional to  $1/v$ .  
The charged particles moving through matter transfer their  
energy preferentially to those electrons that are closer to their  
tracks. The farther an electron

is from the track of the alpha particle, the smaller the impulse it can receive, and hence the smaller  
the energy that can be transferred to it. If that energy is just less  
than the amount required to raise a K-electron to a higher energy level,  
then no energy will be lost to a K-electron at that distance. Some

stint farther any losses to Laloctron will become impossible, and 20

ote The more tightly bound the atomic electrons (Ju:

the atomic number), the shorter will be

the higher

10 "cutoff" distances end

the smaller the rate of energy loss. Consequently the rate of energy

loss is expected to show some dependence on atomic number, being smaller

at high atomic numbers.

---Page Break---

At velocities approaching the velocity of Light, the  $1/r^2$  dependence

is modified by a relativistic factor of

effects The relativistic contraction

of the electric field of the moving alpha particle makes

possible to lose:

Fate of energy loss in ion

at greater distances, and in consequence the

losses slowly at very high energies.

The stopping power is defined as the energy lost by the heavy charged particle (alpha, in our case) per unit path length in the stopping substance and is given by the expression:

$$S(B) \sim \frac{dE}{dx}$$

where  $E$  is the classical kinetic energy of the particles

Stopping power varies with the energy of the particle and the

range of the particle can be calculated by

Equation

$$R(B) \sim \frac{E^{3/2}}{Z_p^2}$$

where  $R$  is the range and  $Z_p$  is the initial kinetic energy of the particle,

where

Stopping power can be determined approximately by measuring the energy of the particles, which have gone through a certain thickness

For a substance, if the range is known as a function of  $Z$ , the stopping power can be obtained from

a 4

= 6.3)

a s(e)

---Page Break---

3+2 Stomrin: Cross Sections The atomic stopping cross section «

is defined as the energy loss

per unit area normal to the

path of the particle, that is:

4)

there  $62/3 \times 49$  the energy loss per unit path length or Linear stopping power and  $n$  is the number of molecules per unit volume. The stopping cross section is used since it is a quantity independent of pressure or temperature,

If the distance traveled by the particle is held constant and the number of atoms along &

Path of the particle becomes the variable

Therefore in which the energy loss depends, the above equation should be written in this form

$$e = -1428 \times 10^{-8} (6.5)$$

In both equations,  $E_0$  represents the energy of the incident particle on the material and  $E$  represents the reduced energy after penetrating  $R$  cm of material.  $n$  is the number of molecules per unit area normal to the path of the particle.

An expression can be derived for  $n$  in terms of the variables

measured in this investigation by the use of the Law of Avogadro, Boyle and Charles with the following relationship:

ok



$a = \rho h \quad (3.8)$

---Page Break---

10

here  $\rho$  is the density of 2

pping uateriel,  $W$  te the atonte

weight and  $A$  ts the Avogadro?s number, Also, ve ean deduct fros  
the atonte theory:

$R_g = m_z \quad (0.7)$

share ng 40 the electronic density and  $Z$  10 the atonte number of the  
substance. Froa Bq. (3-2) ve follove

Row  $\epsilon \quad (1/2) = f \quad (1/n^9) \quad (3.8)$

But ng 1s directly proportional to the density of the materials, 20

?the rengo  $R$  vill be proportional 1). Fron the general law of gases,

$P_Y = H_{ag} \quad G09)$

whore  $M$  ie the molecular naos,  $R_y$  {0 the gas constant,  $P$  4 the pres  
sure and  $T$  Le the temperature in the valune  $V$ . Rearranging the last

expression we get

\*

Equation (10)

From Eq. (3-10) we can see that  $(P/T)$  is directly proportional to  $n$   
and  $n$  is inversely proportional to the range  $R$ . Written in the  
other way,

Equation (11)

---Page Break---

"

Rearranging the last expression,

13)

Under the usual standard conditions and

Equation (6.13)

Inserting Eq. (3.12) into Eq. (3.45) the average stopping power

can be calculated from the collected experimental data using the

following formula:

Eq. (3.14)

where

is the mean residual energy of the alpha particle after traversing the distance  $d$  with a pressure  $P$  in the chamber and the mean residual energy  $E_a$  is the decreased energy of the particle after reversing the chamber at the pressure  $P$  ( $P = P_1 = P_2 = P_3 = P_4$ )  
Eath's Theory. (See reference 14) Bohr's theoretical treatment of the energy loss based on Born's approximation, applied to the collision between the heavy (alpha) particle and the atomic electrons

In this theory the differential cross section for the process in which

the alpha particle transfers a given amount of energy to the atomic

electrons is given by the square of the matrix element:

of the Coulomb

interaction between appropriate initial and final states, Plane waves

are used for the investigations of the

incident and scattered alpha

Particle, the kinetic energies being  $E$  and  $E'$  (S-E), respectively.

---Page Break---

2

The condition of the atom is described initially by the unperturbed atomic wave function:

for the ground state and finally by the wave

function for one of the excited states!

multiplying the result

tion for a given energy loss by the =

energy lost and owing over all

Possibilities gives the final expression for the average energy lost per centimeter of path.

Validity of the Born approximation requires that the amplitude of the wave scattered by the field of the atomic electron shall be small compared to the amplitude of the undisturbed incident wave. As a well known criterion for this is that:

$z_0 v \gg \frac{1}{4\pi\epsilon_0} \frac{q e^2}{\hbar v}$  (3.15)

where

where  $z_0$  and  $v$  are the charge and velocity of the primary particle, respectively and  $\frac{1}{4\pi\epsilon_0} \frac{q e^2}{\hbar v}$  the Planck

constant. This condition is well

satisfied for large velocities and small charge of the incident particle.

Equation (3.15) is also, essentially the condition for the

particle to have its velocity changes when  $\frac{1}{4\pi\epsilon_0} \frac{q e^2}{\hbar v}$  is not neglected,

the particle tending to capture electrons. The evaluation of the

stopping power follows from the velocity of the incident

Particle not only fulfills Eq. (2-15) but also demonstrates, Large  
compared with the velocities of the electrons within the atom, 4.  
an

Be Eq (3.16)

where  $E$  is the energy

the incident particle, 243 the definition,

---Page Break---

B

potential of the electron, and  $M$  is the mass of the incident  
particle and the electron, respectively.

Under those conditions and for nonrelativistic velocities the  
path or "topping power der

2

-a- (see G17)

$$B = \frac{4\pi}{3} n_e Z^2 \left( \frac{Z}{137} \right)^2 \ln \left( \frac{2m_e v}{Z} \right)$$

average energy loss per centimeter

with

where  $v$  is the velocity and  $Z$  the charge of the incident particle,  $n_e$  the number of atoms per cubic centimeter of the material, the nuclear charge,  $I$  the average excitation potential of the atom, and the dimensionless term  $B$  the "stopping number":

Logarithmic term  $B$  the "stopping number",

For relativistic velocities of the incident particle  $\beta \approx 1$

show by Bethe that:

$$B \approx \frac{2}{3} \ln \left( \frac{2m_e v}{Z} \right) \quad (4-22) = 2 \ln \left( \frac{2m_e v}{Z} \right)$$

where  $\beta = v/c$  and  $c$  is the velocity of Light. Although Equations (3.17) and (3.18) were derived for simple "hydrogenlike" atoms, it can be applied to other absorbers by adjusting  $I$ . The value of  $I$  is determined from known experimental data, from which it is found that  $I$  (given in electron volts) is related to  $Z$  by:

loss for  $Z < 0$

19)

loss for  $Z > 30$

---Page Break---

%

From Equation (3.17) the following relationships are evident:

- (a) Stopping power is proportional to electron density of the medium ( $WZ$ ),
- (b) Stopping power is proportional to the square of the charge of the incident particles

Given the stopping power for protons

is  $1/4$  that for alpha particles having the same velocity.

(Velocities are the same for protons having  $1/4$  the energy of alpha particles),

- (c) Stopping power increases with decreasing particle velocity.



---Page Break---

6

A Mock diagram of the detection

system used in this project is given in figure A schematic

of the PND was used to locate the pr

cedure in the chambers The procedure

was within the center as surrounded with @ ceramic canister, Leake

#20 of stopping; gas into the chamber was controlled by a manifold

A 49 alpha source since was placed opposite a scintillator detector

in the gas chamber:

?The number of molecules along the path of the  
Particle which provide the energy loss mechanism ( Coulombic inter-  
actions with the electrons) we varied by changing the gas pressure  
in the chamber:

The residual energy of the alpha particles after  
?traversing the separation distance between the source and detector,  
at a known gas temperature and pressure, was determined from the alpha  
spectra obtained from an analyzer system: consisting of a silicon  
detector (Si) detector, a bias power supply, a miller, a preamplifier, an  
amplifier and a multichannel pulse height analyzer,

4 corners

?schematic of the laboratory setup is shown in figure

42 and a corner

?for close-up can be seen on Figure 43. ?The alpha

Particle source used in this experiment was a calibrated 0.1 microcurie

How is the source, we had a negligible self-absorption and a minimum

sm bactocattorin:, Tats alka source docays by enitting a 5.477-ler

J vas assumed

a we SALSBAIOe (13.6

---Page Break---

6

fem] [oy



-

vecoum

pomp preamp

detector [omelitier

Ly collimators |

Tt?adjustms | height

cylinder | red

chamb

Figure Yo 441: Schenatic dingran of the experizental arrangenents

---Page Break---

Figure 4.2 1 Experizental eet up for  
stopping power and range weesuresents.

---Page Break---

Figure 4.3 +

of the experinental

set ups

---Page Break---

percent) on  $5.37 \text{ AU}$  (1-4 percent). The detector used was of the

sensitive area of  $1 \text{ cm}^2$  and a

2 level of approximately 200

angstroms ( $\text{\AA}$ ) on the insulator dielectric layer on the gold electrode on

silicon nitride barrier type with

typical insulator thickness of

the surface. A range of  $41 \text{ to } 42 \text{ nm}$  dead layer was alpha particle

9 linear stopping power of gold

for alpha particles ( $\text{MeV}$ ) by the dead layer thickness. The

graph showed that  $\approx 5.5, 140 \text{ and } 0.42 \text{ eV}$  alpha particle incident on the

tele energy was constructed using

detector suffered an energy loss in penetrating the dead layer of 10,  
15 and @ Kev, respect

aly, Thus, the detector was essentially free  
of complicating "Nwindow" effects. It has been shown that the res  
ponse of solid state detectors are Linear for alpha particles over  
the energy range of our experiment.

The source and the detector were mounted in line within the  
chamber and the separation distance could be varied by means of an  
adjustable rod upon which the source is mounted. Source and detector  
are shown in figure 44. Also, aluminum collimators were placed  
over the alpha source and the detector in order to get a collimated  
beam of alpha particles

4-2 Energy Loss Coefficient= The calibration factor in Kev/channel  
for the energy analyzing system was obtained by means of a pulse  
observing first the zero energy position and then by observing the lo  
cation of the 5.477 MeV alpha peak while the chamber was under vac.



The eclitrets

save and factor obtatned 4e show in figure 4e5e

The calttration factor (7.77 Sev/chental) wma also noagured after

---Page Break---

oo

Figure 4.4 +

Vacmes char

---Page Break---

srommye/aey 114 gon

Pourqo zoyo8y woRRMMTTH \*woRkTwY ABroue ToD Feol ¥ PUY sOETR ABzotO WO

30 suves 4q pourrge ?uesete SupisveN amod Supddoye or OF eam woRIVIATTED # Sry emlts

¥aEWAN T2NNVHD

008 009 ?oor oz

AOW ?AOWIN3 asInd

---Page Break---

08a ceascrments, Yo aigulficant drifts vere

observed in the calorimetric experiment. The temperature of the laboratory was maintained within 1°F of a mean temperature (75°F) so that it was possible to average several sets of energy loss measurements on each gas.

Using the 5.477 MeV alpha particle and a fixed source detector distance, energy spectra at cover different pressures

were taken for each of the gases considered, Information on the

gases used in the experiments is given in Table 4.1, and an example of several energy spectra of alpha particles after traversing 5.4

cases of air at various gas pressures and at 2, °C 4s shown on figure 4.6. Multiplying the calibration factor (Kov/channel) by the channel number of the alpha peak, we found the average peak energy of the alpha particles reaching the detector after traversing through the stopping gas at a given chamber pressure. Using a simple relation derived from Eq. (3-12) we got an expression for calculating the equivalent

distance traversed by the alpha particle at chamber conditions. In

other words,

Eq. (3-12)

$X_{ch} = X_0 \left( \frac{P}{P_0} \right)^{0.75}$

where the subscripts "0" and "ch" mean standard and chamber conditions

respectively, and  $X_0$  is the extrapolation distance between the

alpha source and detector within the chamber. The last expression

---Page Break---

TABLE 41

SOME CHEMICAL AND PHYSICAL PROPERTIES OF GASES USED IN THIS WORK

cas. molecular weight? standard density

Mt (g/mol)  $\rho$  (g/cm<sup>3</sup>)

Ar 39.95 1.789

Ne 20.18 0.901

He 4.00 0.179

N<sub>2</sub> 28.02 1.251

O<sub>2</sub> 32.00 1.429

CO<sub>2</sub> 44.01 1.977

Freon 12 120.91 4.960

Freon 11 137.17 5.810

5

i, neG-a 16.04 0.717

HH

ue

CO<sub>2</sub> 44.01 1.977

BE

2% BGpe 30.07 4.342

¥

EEE

eat 1878

¥

Rg

e383 Beta 4009 1.967

eae

4s Given by Pauling (92), unless otherwise specified.

21 ahvon By toast (15)," ao)

ja): Given by Orvillo-Thonas|

(b): Given by Bont (2).

---Page Break---

\*ogte 48 pue comnsoext 928 emoTsaA 90 To

30 00 795 Buyszwanny so45e COTOF} TOI MAITU Jo wjoeds ABzuUD Tesanas 1977 OIE

WaGWAN 13NNVHD

oor 009 oos oor oo,

wu 19%

aus 612,

wu 61

|

[SH wus 1

3

2

3

9

WIL 2A11 "235 OOP NI 13NNVH2 Yad SINNOD

8

---Page Break---

may be rearranged as follows

«4.23 (hota)

$T_{tp} = 4 F_0$

assuming that the temperature quotient ( $T_o/T_{eh}$ ) is approximately equal to unity.

Experimental energy and chamber pressure data for each of the runs mentioned in Table 4.1 were taken by varying the gas pressures on the chamber at intervals of 1 to 10 centimeters of mercury and the results were plotted in figures 4.7, 4.8 and 4.9. By differentiating graphically these curves, the stopping power parameters were

calculated taking  $x = 1 \text{ m}$ , and an average value for the corresponding energy. In other words,  $A \approx \frac{E_2 - E_1}{x}$ . Finally, the stopping power ( $\sim \frac{dE}{dx}$ ) was plotted against  $E$  for each gas. The results are given in figures 410 to 413e.

Ranges of alpha particles with energies between 2.60 and 5.47 MeV were determined simultaneously with the stopping power measurements by counting the number of alpha particles of a certain energy reaching the detector in a fixed period of time at a given gas pressure. In order to decrease alpha energies, very thin metal foils (of 0.00004" and 0.000125" to 0.00015" and 0.0002" thick) were placed over the 5.7 MeV alpha source and the energy peaks coming from the decreased energy particles show acceptable broadening averaging. FeO, He,  $^{23}\text{Th}$ ,  $^{232}\text{Th}$ ,  $^{235}\text{U}$ ,  $^{238}\text{U}$  and  $^{244}\text{Pu}$  were obtained using a visual and combined metal foils. Air and hydrocarbon gases were

---Page Break---

at  $^{\circ}\text{C}$  using a  $^{238}\text{U}$  source. The stopping power was measured using a silicon detector. The results are given in figures 410 to 413e.

Figure 410: Stopping power of air for alpha particles.

Z

$\circ$  ? t

new OURN WHETY

---Page Break---

og puw % é2y ty soy osnos0ad v3 zoqueyo om 04 Fuypuodsoss0o +458y

ted upezt09 v Superman, soy ze soToyzrod MTU | Ry ARIL7\*G 30 Azo Dupuyeaey t gry emits

?suis ?Pg 04 Guspuodsossos \*dtsy

i e zg 1

"ROW ?AOWINA VHTY

---Page Break---

qirtoont Pee tn

ary soy omssord cv8 seqavys om 04 luypuedeozros ?Sy yyod upeyTos

© Buyssoaen aeaze soToTaaed MUTE | UY amieU7"S 30 KAsouE DpuyouOY + Gry OmdEE

ssw ?4% 04 Guypuedsossos ?4i8y.

© z 1

"ASW ?AOUINA VHT



---Page Break---

5

u

?Fafo pum ky tape vy soporiswd walye Jo s9oy ABzeHy t OLY eins

?NOW ?9U3N3 WHET

t t f i

a

?un /ROW'AIMOd ONIddOLS

---Page Break---

?new ?ROWINA VHaTY

+ °

?w2/how ?WIMod oNlaois

---Page Break---

?

§

09 puw Zo ?2y uy copoyssed ytre Jo ssoy sour t ziry oman

?ROW ?OWING VHdIV

su2/A9W ?YIMOd ONIddOIS

---Page Break---

rh moons pum ofa Zax ?Rie wy eoToTa Tad melTe Jo ssoT Sew 1 Chey oT

?Now ?kON3NT VHGIV

Z

---Page Break---

B

chosen for this part of the axperinents end tho resulte aro tabulated

98 Table 4.2 and plotted on Figures 4.14 to 020. ditional range

date for other gases auch as O25 M2, 002, A2, Erp, MgO and Freonet4

48 presented on Figures 4.21 and 4.22 for 5.47-View alpha particles  
Range curves were calculated by normalizing the total counts using  
the ratio of total counts and maximum total counts recorded by the  
scalar for the same period of time. Table A.1 (see Appendix) is a record of the

things of the instruments used on the experimental

Measurements.

4-3 Gross Section Calculations.

and hydrocarbon gases (C<sub>1</sub>, C<sub>2</sub>, C<sub>3</sub>, C<sub>4</sub>, C<sub>5</sub>, C<sub>6</sub>, C<sub>7</sub>, C<sub>8</sub>, C<sub>9</sub>, C<sub>10</sub>) were calculated

by means of Equation (3.14) and the following constant values

Avogadro's number (A):  $6.024 \times 10^{23}$  molecules/mole

Standard Temperature (T): 273°K

Molecular cross sections for air

Standard Pressure (P): 760 mm of mercury.

Source to detector distance (a): 5.4 cm.

Molecular weight and standard density values for each gas to

be inserted on Byström (3.14) are given in Table 4.1. Cross Sections

for air, nitrogen, ethane, propane, ethylene and propylene for 0.3 =

5.0 MeV alpha particles are tabulated on Table 4.3, and curves of

molecular cross sections ( $10^{-24}$  cm<sup>2</sup>/atom versus alpha energy

(eV) were drawn in Figures 4.23 and 4.24.

---Page Break---

REMEDY ROLAT

PARTICLE

BERT,

MeV

5047

5.04

48h,

423

4.06

3653

3638

227

2.60

1643

3.89

3653

3432

2B

2.63

2.10

141

TABLE 4.2

a

427

3674

3652

242

1.92

178

wn

SIP OF AIR AID HYDROGAR

FOR ALPHA FARTTOLEs®

Rue, cus

a, alg

2.96

2.38 -

213 2.00

1631642.

1531435

W19 1.06

0.69 O68.

\* These values are plotted in Figure 4.2s

GASES

ose

2.05

157

4.51

wit

1.03

0.92

0.62,

ora

0.70

0613

---Page Break---

38

soandy vary uy erFOR TL Pim nO UTI

uw (Aixeuo ofsoupy TerATEy A

seme Sao nee soars 07m

edsowg aoasU soToFIZOd WydTS s0y

sororszed sutte ns 01% 10J coamo oRMy sYL+y omigg

sm Py on

Et 1 i

Surpuodsosso ?Aisy.

(s4un02 -xou/ssun0>) Olayy

---Page Break---

sa0y wy teonuA Ho woE somo amy uO complE

\*(Ya0) OweNOK UF OTROS KL puW nO UIIA ArOK Susesdsony amaze coTorired wate TOF

Pew (Cisous OFOUET THRAFEE aml L7°6) GoTo wylT | ey OU 205 CoamD OOMY 4 SLH7

OmdEE



San EE SUEDE EPEAT SE URL eee

Fun Py 01 Buypuodscison ?Airy t

ou syuno>) DLAVe

(snunes \*

---Page Break---

7

os

He

| t

say UF foonpes Msous wooH soAm2 OWN Wo EoD TE

\*Ci®o) oumna wy stTos FL pu ng UT fren BuyecHdsen aoyse soToTAcod wydTe Joy

uw (A208 OpROUTy THFRFUT Ami LY°S) SOTSTA TEL BUITR | RY om JO soamo euMy 1 GLY

omITE

Su Jo sw ?aunssaue waawvH> ?V4

ov co or

ssa ?Py 04 Surpuodse.

oO

(24un02 -xow)sune>) O1vy

---Page Break---

< sami wy teougeA ABzau9 wee seams omy WO CORTE

+@afo) eundorg wy sTTos WE PHO TO UTI Lx0n Jupoodoon aoqJv coToTyTed wdTE TOF

Paw (Gisoue OTOUFE THFATUT a0H LY°6) FOTOHAIOL MRT | Aa omy op coaMD oUNY t LLY

amdys

SH Jo sur ?aunssaya uzewvH ?Ys

os or or ox a

jun) OUVA |?

(Sunes ous

---Page Break---

»

Cio) ovovtina wy sos FL ee

Pov (£35000 STOUT THFETU AOL)

os

wHeD

osu ?1G

poo

of Buspuodsesso> ?45.

aire 03

woamo odimy + ghry oma

or

ese

2 2

(squno> -xow /syuno2) Dive

=

=

---Page Break---

g

\*(i9) oworkoas up oTFOE FL PED ny

os

conte Kin even cone of vo eomdRE

9a SupeceLIONy xOAJ0 COTATI mAIPU GOs

uD (4Bx0us OTROUTL TERA FUT AM TOTATE WSTY | HY Om OF Som eum IGLe7 omiry

8H fo sup ?gunssaua wagwvH> ?YY

ov or ot

a

ouvea ?

8

2

(s4und2 ?xu /ssunoy

i

su ?04 Buipuodsesso> ?4ity

---Page Break---

RANGE, ems

1

i 2

ALPHA ENERGY, MeV

Figure 4.09 + Ransom:

crete

tone

a

Calg

Coty

CoH,

CoHs

5 6

ba particles

underd condi

---Page Break---

(ou paw % ey

24) woees ovmszouy Texwaer uy (40H L7\*5) woTeTAEM MYSTY j,\_t7 sO} Conan BDU 1 1Z~y

emoqs

Sy 4o sus ??aunsszus usewvHd ?Pa

9 os ov ge L

oz

ot

° sun "y §, buypuodsosso> ?dasy ©

> Youve

(s4tin09 oui /syuno:

---Page Break---

7h wo0ag pas a

oo wr (a5 LYS) seTONIOd wHET® 7 403 Boasmo eBUWY t ze+7 OTE

6H fo sus ?sunssaua wagwrno ?Ps

os op oe Sr m

stuns) O1LWs

to

(squno> -xoui |



---Page Break---

TABLE 4.3

MOLBOULAR STOPPING CROSS SDOTIONS OF AR #11D HYDROCARBOW GASES  
FOR ALPHA PARTICLSS\*

anh, tay tae OCOe

0.3 - - 13.5 196 235 26.0

Oh - - 16.0 213 28.5 27.0

05 55 - 1762 225 21 215

ob aks mse

er ee

12 9 1S 265 16.0 15S 22.6

166 5S 663 1064 1267 We Wd

2.0 8.0 55 cy) 10.6 1267 1406

28 55 AS 7600 8.0 10-1 Wed

32 ?7 rd 72 93 1065

---Page Break---

CONT. TABLE 443

nem

oo,

vat

at

36

ae

40

42

bs

b6

be

50

\* Those values are plotted in Figures 4.23 and 4e2be

am

be

403

bet

39

28

Be

37

366

35

MOLECULAR, STOPPING CROSS SECTION

(10<sup>-4</sup> eV/aolecule/en\*)

my

309

308

36

35

33

32

3A

3.0

29

Cai,

6.2

6.0

58

566

Soh,

50

5.0

49

he

ale

69

67

665

63,

Gt

569

BT

56

55

ORG

89

8.6

83

8.0

18

16

Th

13

ms

Cig

10.0

7

Qh

9.0

87

8.6

8.4,

8.0

---Page Break---

a4

8

MOLECULAR CROSS SECTION,  $10^{-24}$  cm<sup>2</sup>/molecul

3

1 2 3 4 5

ALPHA ENERGY, MeV

Figure 4.23 : Molecular stopping cross sections of air, O<sub>2</sub> and

CO<sub>2</sub> for alpha particle:

---Page Break---

air

3

8

8

MOLECULAR CROSS SECTION, 16!ev/molecule/cm<sup>2</sup>

3

ALPHA ENERGY, Mev

Figure 4.24 : Molecular stopping cross sections of Gly, C<sub>2</sub>H<sub>4</sub>, and

O<sub>2</sub> for alpha particles. ?ar Cal,

---Page Break---

## CHAPTER

### RESULTS 21D corrections

The accuracy of the stopping power and oscillator cross sections calculated in this investigation depends on the following quantities:

**Je- Gaa Pressure:** The difference between the two columns of the Mercury Barometer could be read to an accuracy of  $\pm 0.5$  mm of mercury.

Thus, the error in  $P$  ranged from 0.4 percent for the largest pressure intervals, to 6 percent for the smallest pressure intervals. The pressure intervals ranged from 113 mm of mercury at the highest energies to 8 mm of mercury at the lowest energies.

2

**Gas Temperature:** A mean temperature of  $21^\circ\text{C}$  was used in calculations and the temperature of the laboratory varied  $\pm 1^\circ\text{C}$  of this mean. Therefore, the fractional error was approximately 4.0 percent.

**Separation Distances of source and detector:** The error in the measurement of the separation distance was of the order of 0.8 percent.

**Residual Energy of the alpha particle!**

This error de the

most important in these experiments and the hardest to estimate, However, the estimate of this error

is reasonable from the spread of

the experimental data obtained from the energy loss spectra The uncertainty in estimating the peak channel of the energy spectra was

Rayon power of error in determining the average

At high residual energies, the resolution of the energy

system was very good and the error:

Locating of the energy



---Page Break---

49

spectra were noted: » straggling of the particles caused the energy spectra to broaden and the peak could not

be located with as much accuracy, in this

» the error was estimated to be 2 percent, or 0.23 percent at higher energies and 5.5 percent at the lower residual energies.

The probable error of the calculated values based on the experimental data was determined by taking the square root of the sum of the squares of the maximum fractional error in each experiment.

At higher residual energies!

99 the most probable error (mean standard deviation) found was of the order of 4 ~ 5 percent and 9 ~ 10 percent for lower residual energies respectively.

The value proposed by the author on the basis

of the investigation

for the stopping power of 13 cases are given on figures 4.10 to 4.13 as a function of alpha energies; also, for the molecular cross sections of air and hydrocarbon gas (Gly, Calg, C26, CHR, Osis)

which are tabulated on Table 4.1. Those energy velocity values were obtained from smooth curves drawn through the experimental data points and probable errors (figures 4.23 and 4.24) as a function of alpha energy; this is also true for the stopping power and range curves.

The method used in this investigation gives values which are in excellent agreement with those found when the gas pressure within the

chamber is constant and the scintillator distance is varied. 4a

the energy region where a separation is possible, Table 5.1 is an

example of comparison of molecular energy:

15)

3 section calculations given

by the authors find this work showing a good correlation within

---Page Break---

U6

670

een

on

ort

ru

cron

oe

ove

ore

oe

zo bg

ah ug

ek ae

oe

g

rot

eu

sa

ote teh

ee sest

oe Stak

vor

v

Ms

z9

ow.

we

\$6

oro.

6%

Ea

be

29

ber

es

or

29

ze

oe

ba

8

re

66

\$6

se

ore

se

or

at

ot

mm

zn

on

80

90

RAT Wee ay Te

8

uly

execu

Sioa

way

SRUQUVILT? MO GNOOE VAVO LIA SYNGE NOLLOaS sooND UVINLETOH go NOsTEVAHOD

bs save,

---Page Break---

51

ow cL BT wr 6 He se oe ors

oe LS sae re ue ore or

Be 9S ss HE ue ore or ov

Te ST | AE ay FE

) SIETDEC Cro) annie ary

gio/OTAOTOH/AW OL \*NOLICES SOND WrUNATION

---Page Break---

the

= 5:0 sav alla energy reston for soue hydrocarbon gases and

air, This vo:

could be extended indefinitely for additional data

outer cross axis

isn calculations using the information on inorganic

gases (type, Copy 0, etc) given in this work or discussing as

many foreign and organic research and industrial gases as desired,

In the same way, comparisons with theoretical calculations using the

Bethe's equation (Eq. 3-17) agree when the residual alpha energy is

higher than 1 eV. (3

Figure 5.1).

In the radiolysis of a gaseous system it is possible to measure

the number of ion pairs formed by absorption of radiation. The yield



of a given rea:

on can therefore be expressed in terms of the number  
of molecules

30d per ion pair formed. The ratio  $N_{IP}/M$  is commonly  
referred to as the ion pair yield, where  $M$  is the number of molecules:

changed per ion pair and  $N_{IP}$  is the number of ion pairs formed.  $M$  may  
denote a species of a given kind disappearing or being produced,  
and in

so of cases,  $M$  may be expressed either as the total  
number of ions or as the number of molecules of a given kind. The latter  
calculation is based on knowledge of stopping powers and the energy  
required to produce an ion pair in the individual gases. Related re-  
search could be developed simultaneously. Like gas radiolysis, deter-

ination of chentel roactin cechanigns, on ylold seasuronanta, etoy  
using tho values given in this vork for inongente and hydrocarbon  
azo:

---Page Break---

Alpha Enorey, Me

Figure 5.1 Ranzeenor:

3

jethe

slow alpha partiele

---Page Break---

cf

he

om

10.-

BTBLTOORAPEY

ALLIS

Seis and W:RSEAN, \$.D. The Stopping Grose Sections

for Protons,  $O^{+1}$  to 6.0 Mev} Rev. Mode Phys 25, 779 (1953).

BRIT, HAs: Chentel Structures of Incrsante Gases} Je

Baue. 37, 616 (1960).

BENE, HA,

The Rango-hergy Relations Zor Slow Alpha Parti=

les and Protons in Air} Rev. Mods 2p 213 (1950).

+ GHD, Weis and POWERS, D.t Stopping Cross Sections

for Alpha Particles from 0.3 to 2.0 Movs Phys. Reve

BB, 3625 (1971).

COLARD, Js and GAL, J+ Silicon SurtacoZarrier Type Sentcon=

ductor Detectors; lucl. Instrim, Moths 16, 195 (1962)

BOURLAND, Pe

DEARIALEY, G.: Semiconductor Nuclear Particle Detectors; National  
Academy of Sciences ~ National Research Council Publication  
871 (Edited by J.W.T. Dabs and FeJe Walter), ppd?  
(1961).

JESSE, W.P. and SADADSEIS, J,

The Range-energy Curves for Alpha  
Particles and Protons; Phys. Rev. 72, 1 (1950).

KERR, G.D., HATRR, L.

My, UNDERWOOD, MN, and WALTER,

4 Mole

cular Stopping Cross Sections of Air, Ygy Kr, 002 and Cl,  
for Alpha Particles; Health Phys. 12, 1475 (1956)«

EMD, 5.0.1 Radiation Chemistry of Gaseous Reinhold Publishing

Corp. Yow Fors, Ne. pps 123 (496%)«

ORVILLE, WSs and THOHSS, Tet sade Spaeteye 3, \$89 (1955)e

---Page Break---

55

Wee PALER, BBS Th

Stopping Power of trogen and Rydrocarbon

Vapours fer Alpha Particloy over the Energy Range 1 = 8 Mevy

Proce Fhys. Soc. 87, 681 (1966).

Y= PAULING, tt The Chenteal Ponds Comal? University Press, Ttheca

MWeEs pe 58 (1967)

Bee ROTONDT, 5.

Energy Loss of Alpha Particles in Matter; Radiation  
Base 33, 1 (1968).

Vern SEGRE, B (Bitter): Experimental Nuclear Physics vol. 1, part II,  
pp. 166 Wiley (1953).

Wen MALAH, Pad.

Shory Lose of Alpha Particles, Protons and Ele  
trons in Matter; Ph. D. Thesis, University of North Carolina,  
Chapel Hill, N.C. (1962).

WEAST, R.C. (Baird): Handbook of Compton and Photoelectric  
Effect; McGraw-Hill Book Co., Cleveland (1966).

---Page Break---

---Page Break---

57

TABLE At

INSTRE-STATION CAUIPMENT SETTDUS FOR EXPERDGNTS

Joe DETECTOR BISS SUPPLY

Ortec Model 423

Gate Ar 50 volta.

mR

Ortee Motel 719

Ortec Nodel 486

Coarse goin: 16 volts

Pino gaint 6.5 volts

Window: 10 volts



Lower Level: 0452 volts

tsar settings as per operating instructions.

dew Scatsn

Ortec Yodel 484

Trroshalas 0.4

Other settings as per operating instructions.

ARALOG TO DISTTAL co:TSATER

Mucloar Data

Yaro Fino: 0.10 volts

Tero Coarse: 10.0 volts.

---Page Break---

58

TUBE At

Coaverston Gains 102% che

Upper Level Diserisinator: 10 volte,

Lower Lovel Discriminator 0.49 valte

other sett

as per operating instructions.

Gem MASTER CONTROL

Nuclear Data

Reset time: 400 seconds.

Other settings as per operating instructions.

?Tom READ-TA/OUT DISPLAY

Master Data

Settings as per operating instructions.

---Page Break---

TABLE A.2

ROE AND EC MOE Losses measured at 5.477100 aT anrma

PARTICLES DN ATR

ona of Hg = MeV

os tm tam oak

10.5 1643 0.962 0.746 622 ABS

re

1669 1603 0.939 14200 557 432

3067 1644 0.963 20181 423 3629

3563, 1623 0.950 2.508 373 29

3903 1621 9.949 2792 33 2650

427 1627 06953 3.033 2 210

M7 1629 06954 36318 wz 1.00

5103 1549 06907 3a6h4 13 0.87

5360 1085 0.635 36765 R O71

5he8 36 0.002 3.893 2 0645

---Page Break---

TABLE A.3

AOSE gm Remar toss mr=noewes, pava FoR 5.477N6v ax?41 goon

PARTICLES TY i,

caasaR ?tora, CRANE = QUTVALENT

sta, comms RAN Sure ?BIEGY,

cas of Hg = We

nk 16 1.000 0.028 eg 5h

5.0. 163804954 0.355 gy 52

3.0 170.973 0.639 a5 5.01

1364 169% 04937 o.930 5 478

1k 16690972 1.236585 454,

22 1639 0.955 1.506 56 432

25eh 1628 0.954 14804 502 4.05

32h 1603 0.934 2502462 3659

36.2 1686 0.982 572 yap 3.32

ed, 1580 0.92 aur gre 2.89,

09 105 0.935 3.332 2645

5204 WI 0435 371g 1093

5307 71 0.040 3.815 at 1663

Boor a.843 - -

8 04004 3.986 179 1.39

104000 4227434 1.02

10,000 kez70 8 0.65,

---Page Break---

TABLE AL

SANOZ AND DWERGY 10SS EXPEROENTAL DATA FoR 5.47-¥EV an arena

PARTICLES Di Cig

cmeen ora, can, EQUIVALENT

rassuss, courts, «ATO Inte za, ?ENERO,

eas of Hg os Mav

oe 1663 0.999 0,056 698 Seka

36 1602 0495004255662 5h

10 1636 06974 a ho82

907 1685 1.000 ons 586 455

127 1583 0.999 got 542 bez

1564 1584 0.940 14093505 3692

18.6 168 06954 14320 ass 355

216 1518 0.936 14533408 3.17

2b 20.956 6353 2.1

2.0 Kb 0.976 4.968 a7 245

316 1567 0918224319 148

3h8 ee 0.408 KB 0.56

3504 58 00m 25133 0433

36.0 © 0,000 24556 5 0.19

---Page Break---

TABLE a5

RANGE aD ENERGY Loss EXPERDENTAL para POR 5.47-NEV An2%1 ana

cE

PRessuas,

esa of Be

1.0

3.0

25.5

26.3

?ror,

cours

1613

1631

163

1643

15%

1604

1565

1952

us?

1430

546

202

1

RATIO

0.959

1.000

06965

oom

0.949

0.954

0.930

04923

ong

04850

06324

0.120

0.000

PARTICLES TX O3iig

ee

o.071

0.355

0.540

ov7e8

1.072

1s

148

46541

1.669

ame

1.769

nen

1.850

cae

BER

693

623



573

503,

an

35h

224

uo

2

65

8

BqUIvALENT

?Bier,

Mev

5638

8h,

eks

3490

3.20

2.75

2.20

wn

1.08

0.95

ont

0450

oont

---Page Break---

TABLE 6,

RKRS AND BRmey 109s exrsenc:

PAXTIOI

ral DATA FOR 5.4.77 an%47 popag

Ti Cait,

Ri TOME cum, agora

cE, COUNTS RATION yp, © NRE ?BIEAGT

% Ea Mev

07 1699 0.990 0.049 m2 5eh5

29 WHS 4.000 0.206 69 5.27

509 1705 049% 0.419 as 5.04

89 Wet 04980 0.632 616 4

19 1682 0.900 0,045 sz 452

Med 1604 04935 1.058 33 heat

12.0 168104980 1.207 504 3.98

21.0 W641 04956 14492 465 3662

200 1680 0.989 44705 re) 3.28

21.0 Wis 0.960 1.918 380 2.95

30.0 15610910 291 325 252

33.2 162004980 24358 267 2.07

36.0 1503 0.876557 206 1.60

3700 15K5 0,900 2.608 m4 1635

3807 150604878 a9 190 4.01

4004 1254 O71 2.849 88 0.68

4009 200 0.116 2.906 5h ona

43 0.012 2.934 52 040

au? 10,0005 2.962 x» 0.23

---Page Break---

LE 7

RAWE AKD EER 1089 DOAROGTAL DATA FOR 5-4 7-30V an41 arpea

canaER

FRESSIRE,

ens of Ez

0.8

38

68

38

128

15.9

18.9

2109

25.0

21.0

2.8

Bh

28.8

RATIO

1.000

0.9h4

0.959

0.4985

0.976

0.981

0.962

04943

04932

ow7es

0.207

0.002

0.000

PARTICLES Dt Cg

Tatpy

ee

0.056

0.270

04483

04696

04909

1.129

1.343

16555

177%

1918

1.915

2.017

2.046

cue,

HOGER

6%

653

00

550

494,

431

360,

186

105

R

32

SquaaT

'RIERGY,

Mev

Soh

5.07

467

beet

3683

3634,

2.79

2.28

ohh

0.81

0655

0624,

0615

---Page Break---

SE AND

oO?

Det

67

10.1

1363

1669

20.8

25.8

324

Bue?

38

35.7

36.1

369

TUBULE 8

Y L083 EAPERDSEMAL ATA FoR 5.4770 aT arena

TOTAL

cours

16

1706

1670

167

1653

1685

1619

1565

1592

1532

1399

s92

153



10

2

PARTICLES TH 0

RATIO

0.991

012

0.991

04993

0.981

1-000

0.960

0.950

0.945

0.910

04830

0.707

0.090

0.005

0.000

94050

04220

06476

ont

0.945

1.200

wart

1.833,

1.925

2.280

2.420

2.472

2.536

2.565

24621

REER

699

667

62

sm

530

45

27

on

169

109

92

&

49

BRUIvalrr

?BIERGY,

Xe¥

50h

5018

4083

hel

an

3668

3624,

2.67

2.10

4.31

0684,

ont

0652

0.38

0.20

65

---Page Break---

1528

1454

120

TABLE A.9

ISLES TW 02

RATIO

1.007

0.995

0.957

04959

o.ous

1.009

0.958

ome

0.960

04970

068%

0.901

04904

0.860

0.275

0.071

0,000

cua,

RBI

703

232

133

DATA FOR 5.47=RCV An°%47 aupna

EQUEVALST

?BERRY,

Nov

Sekb

53h

5.06

4480

bokh

beth,

38h,

308

3605

2.70

2.16

1.80

1.07

0.87

0656

0.42

0.26

---Page Break---

88

m8

8

2.8

32.8

lh

41.0

45.0

Mae

506

52.0

5208

5309

5563

5548

56.2

1697

1758

1688

1652

1588

1624

1567

1651

1607

1626

16%

1560

1513

173

ea)

1939

389

?TABLE a-10

RATIO

1.000

1.035

06994

0.973

0.935



0.956

04923

0.973

0.947

04964

0.951

0.919

0.891

0.9%

on8e3

0.789

6dr

o.217

0,010

0.001

Tater

0.056

0.341

0.625

1.051

1.477

1.762

2.046

24330

2.657

2.913

2197

36467

34595

34694

3.751

3.829

3.826

3.929

3.964

3.993

msm

705

or

67

590

53h,

496

458

412

361

3h

203

m

138

1%

a2

n

56

ar

S2ITAL DATA FOR 5047-4 ant apna

PQUIVALET

BHEROY,

Mav

Soll?

5622

9h,

458

hath

3.85

3655

3.20

2.80

2043

2,02

157

1033

1.07

O97

0.76

0463

0655

0643

0.36

67

---Page Break---

5545

5643

5965

599

TaBte At

EPrEnm

PARTTSSISS

Tra

Xstor

o.170

06454

o.n7

1.001

146342

1.612

1.911

2.223

2.479

2.263

36133

36431

3.701

e255

crane,

mma

690

659

598

559

523

182

125

102

52

TAL DATA FOR 5.,77%MEV An? urna

EQUIVALENT

?BERGE,

Mev

5636

512,

487

dob

beh,

4.06

31h

Bet

3et2

2663

2.20

1.87

et

om

0.79

0654,

0440

0.20

---Page Break---

TABLE At?

04998

0.970

1.000

0.980

0.994

0.974,

o.o71

04959

o.g71

0.934

0.907

0.927

0.919

o.922

0.892

04566

0.025

0.000

?a

0.028

0.099

0.227

0.362

06518

06675

nese

1.072

1.321

1.463

16

1748

1918



207

26344

2.499

2.570

2.644

came,

UGE

700

687

669

640

615

588

505

BRR ERS

1st

22

56

SRESTTAL DATA POR S477 aT gpa

PARTICLES TN Op

PQUIVALETT

?SERGY,

av

5h

5633

5.18

4

am

4456

422

3.92

3655

69

---Page Break---

TABLE 2.13

PAIGE 27D EVEGY Loss E@IMETAL para Por 5.47ev An? appmy

PARTICLES Tt Ee

cana

PRESSURE,

cas of Ee

O64

bok

94

Boo

18.9

25

22

31.0

33.0

35.0

3569

3667

3803

rom

conts

1657

1687

1662

1662

1567

Mah

104

conn a 8 ER

RATIO

0.982

1.000

04985

0.985

0.928

0.880

0.832

0.798

06138

0.017

0.004

0.000

(0.000

0.000

0.000

?stor

0.028

04312

0646

0.987

16342

1.669

1.932

2202

26344

26415

26486

2.550

2.607

a1

2.842

cua,

OS

4,

SQUIvALErT

?BERG,

Mev

Solet

5.08

4063

hts

3.67

3.1%

2%,

2.20

1684

?7

1.60

1643

1.30

1404

0463

---Page Break---

8

38

68

10.1

12.9

15.9

18.7

19.8

200k

21.0

RY Less

ror

TABLE etd

SPERD IAL pan

PARTICLES Tr

cours ?Rarro

1679

1683

1645

1656

1630

1625

154k

103

112

0.968

04965

0.917

0.714

01086

(0.000

ron s.c77?v an arpa



693

3662

2.96

217

116

0.70

648

0.25,

n

---Page Break---

---Page Break---